

# q-Euler Lagrange Equation

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## Article history

Received: 03-10-2019

Revised: 04-11-2019

Accepted: 22-11-2019

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**Abstract:** In this work, for  $q \in (0, 1)$ , the q-deformation of Euler-Lagrange equation is studied; we use q-derivative (or Jackson derivative) to deduce a new formula of Euler-Lagrange equation. Some examples are presented.

**Keywords:** q-Derivative, q-Euler-Lagrange Equation

## Introduction

The calculus of variation is a one of the most important division of classical mathematical analysis as regards application. The aim of this study is to supply the reader with a certain minimum of problems covering the basic divisions of classical calculus of variation. In Euler; s methods, the values of the functional see Equation 1 are considered not on arbitrary curves admissible in the given variational problem, but only on polygonal curves composed of a given number  $n$  of straight-line segments with specified abscissas of the vertices. If we have a certain class  $M$  of function  $y(x)$  and each function  $y \in M$  there is associated, by some law, a definite number  $J$ , then we say a functional  $J$  is defined in the class  $M$  and we write  $J = J[y(x)]$ . The class  $M$  of function  $y(x)$  on which the functional  $J[y(x)]$  is called the definition of functional. In mathematics, a q-analog of a theorem identity or expression is a generalization involving a new parameter  $q$  that returns the original theorem, identity or expression in the limit as  $q \rightarrow 1$ , see (Altoum, 2018a; 2018b; Rguigui, 2015a; 2015b; Bangerezako, 2004).

The Euler-Lagrange D.E is the essential equation of variational principle. It is defined by an integral of the form:

$$J = \int f(t, y, y') dt \quad (1)$$

Where:

$$y' = \frac{dy}{dt}$$

Then  $J$  has a stationary value if the Euler-Lagrange differential equation:

$$\frac{\partial f}{\partial y} - \frac{d}{dt} \left( \frac{\partial f}{\partial y'} \right) = 0,$$

is satisfied. If time-derivative notation is replaced instead by space-derivative notation, the equation becomes:

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \left( \frac{\partial f}{\partial y_x} \right) = 0.$$

The Euler-Lagrange differential equation is implemented as Euler equations  $[f, u[x], x]$  in the Wolfram Language package variational methods. In many physical problems, the partial derivative of with

respect to turns out to be 0, in which case a manipulation of the Euler-Lagrange differential equation reduces to the greatly simplified and partially integrated form known as the Beltrami identity:

$$f - y_x \frac{\partial f}{\partial y_x} = c.$$

For three independent variables, the equation generalizes to:

$$\frac{\partial f}{\partial u} - \frac{\partial}{\partial x} \frac{\partial f}{\partial u_x} - \frac{\partial}{\partial y} \frac{\partial f}{\partial u_y} - \frac{\partial}{\partial z} \frac{\partial f}{\partial u_z} = 0$$

This paper organized as follows: Section 1 and 2 introduce the basic concepts of the study. In next section, we present an analogous of the classical Euler Lagrange equation, finally we deduce some theorems.

### Preliminaries

We introduce the q-Derivative, we recall some basic notations used in q-calculus. The natural number n has the following q deformation:

$$[n]_q = 1 + q + q^2 + \dots + q^{n-1}, \text{ with } [0]_q = 0.$$

Occasionally, we shall write  $[\infty]_q$  for the limit of these numbers:  $\frac{1}{(1-q)}$ . The q factorials and q binomial coefficients are defined naturally as:

$$[n]_q! = [1]_q \cdot [2]_q \cdot [3]_q \cdots [n]_q,$$

With:

$$[0]_q = 0.$$

Here is a decent redirection for any individual who knows what the derivative of a simple function is  $f(x)$ . The modern theory of differential and integral calculus began in the 20th century with the works of Newton and Leibniz. As it is well known, see (Altoum *et al.*, 2017), the derivative of a function  $f(x)$  w.r.t the variable  $x$  is by definition:

$$f'(x) = \lim_{h \rightarrow 0} \frac{f(x+h) - f(x)}{h}.$$

Now, let us consider the following expression:

$$f'(x) = \lim_{q \rightarrow 1} \frac{f(qx) - f(x)}{qx - x}.$$

Of course, this is not valid when  $q = 1$  or  $x = 0$  but otherwise this alternative formula is equivalent to the usual derivative. You can convince yourself by writing  $\frac{f(x+(q-1)x) - f(x)}{(q-1)x}$ , the term  $(q-1)x$  playing the role of  $h$ .

At the beginning of the 20th century, F.H. Jackson studied this modified derivative and many of its consequences, see (Rguigui, 2016a; 2016b; Rguigui, 2018a; 2018b). The key concept is the q-derivative operator defined as follows when  $0 < q < 1$ :

$$(D_q f)_{q \rightarrow 1}(x) = \lim_{q \rightarrow 1} \frac{f(qx) - f(x)}{qx - x}.$$

This q-derivative can be applied to functions not containing 0 in their domain of definition. Then it reduces to the ordinary derivative when  $q$  goes to:

$$\lim_{q \rightarrow 1} (D_q f)(x) = f'(x).$$

As an example, we compute the q-derivative of  $x^2 + 2x + 1$ :

$$D_q(x^2 + 2x + 1) = \frac{[(qx)^2 + 2(qx) + 1] - [x^2 + 2x + 1]}{qx - x}.$$

One can easily check that the q-derivative operator is linear:

$$D_q(f + g) = D_q f + D_q g$$

$$(D_q \lambda(f)) = \lambda(D_q(f)),$$

the product rule is slightly modified but it approaches the usual product rule when  $q$  goes to one:

$$(D_q(fg))(x) = f(qx)(D_q g)(x) + (D_q f)(x)g(x).$$

### q-Euler Lagrange Equation

As analogous of the classical Euler Lagrange equation:

$$\frac{\partial f}{\partial y} - \frac{d}{dt} \left( \frac{\partial f}{\partial y'} \right) = 0 \tag{2}$$

Where:

$$y' = \frac{dy}{dt}$$

We introduce the q-Euler Lagrange equation as follows:

$$D_{q,y}f - D_{q,t}(D_{q,y_q'}f) = 0 \tag{3}$$

$$l(t) = \frac{f(t, y(t), y'(t)) - f(t, y(t), y'(qt))}{y'_q(t)(1-q)}$$

Where:

$$f = f(t, y, y'_q)$$

$$y'_q(t) = D_{q,t}y(t) = \frac{y(t) - y(qt)}{t(1-q)} \tag{4}$$

$$D_{q,y}f(t, y(t), y'_q(t)) = \frac{f(t, y(t), y'_q(t)) - f(t, y(qt), y'_q(qt))}{y(t)(1-q)}$$

Then, for  $q \in (0, 1)$  the q-Euler Lagrange Equation (3) is equivalent to:

$$\sum_{k=0}^{n-1} q^k \left\{ \frac{f(q^k t, y(q^k t), y'_q(q^k t)) - f(q^k t, y(q^{k+1} t), y'_q(q^k t))}{y(q^k t)} \right\} \tag{6}$$

$$= \frac{l(t) - l(q^n t)}{t}$$

and:

$$D_{q,y_q'}f = \frac{f(t, y(t), y'_q(t)) - f(t, y(qt), y'_q(qt))}{y'_q(t)(1-q)} \tag{5}$$

*Proof*

Let q-Euler Lagrange equation is given by  $D_{q,y}f - D_{q,t}(D_{q,y_q'}f) = 0$ . Let  $l$  given by:

$$l(t) = \frac{f(t, y(t), y'(t)) - f(t, y(t), y'(qt))}{y'_q(t)(1-q)}$$

From the above discussion, we obtain the following theorem.

**Theorem 3.1**

Let  $l$  given by:

Then, we substitute (4) and (5) in (3), we get:

$$\left( \frac{f(t, y(t), y'_q(t)) - f(t, y(qt), y'_q(qt))}{y(t)(1-q)} \right) - \frac{l(t) - l(qt)}{t(1-q)} = 0$$

$$\left( \frac{f(qt, y(qt), y'_q(qt)) - f(qt, y(q^2t), y'_q(qt))}{y(qt)(1-q)} \right) - \frac{l(qt) - l(q^2t)}{q(1-q)t} = 0$$

$$\vdots$$

$$\left( \frac{f((q^{n-1}t), y(q^{n-1}t), y'_q(q^{n-1}t)) - f((q^{n-1}t), y(q^nt), y'_q(q^{n-1}t))}{y(q^{n-1}t)(1-q)} \right) - \frac{l(q^{n-1}t) - l(q^nt)}{q^{n-1}(1-q)t} = 0$$

Then, we obtain:

$$\left( \frac{f(t, y(t), y'_q(t)) - f(t, y(qt), y'_q(qt))}{y(t)} \right) - \frac{l(t) - l(qt)}{t} = 0$$

$$q \left( \frac{f(qt, y(qt), y'_q(qt)) - f(qt, y(q^2t), y'_q(qt))}{y(qt)} \right) - \frac{l(qt) - l(q^2t)}{t} = 0$$

$$\vdots$$

$$q^{n-1} \left( \frac{f((q^{n-1}t), y(q^{n-1}t), y'_q(q^{n-1}t)) - f((q^{n-1}t), y(q^nt), y'_q(q^{n-1}t))}{y(q^{n-1}t)} \right) - \frac{l(q^{n-1}t) - l(q^nt)}{t} = 0$$

This leads to:

$$\sum_{k=0}^{n-1} q^k \left\{ \frac{f(q^k t, y(q^k t), y'_q(q^k t)) - f(q^k t, y(q^{k+1} t), y'_q(q^k t))}{y(q^k t)} \right\} = \frac{l(t) - l(q^n t)}{t}$$

Which completes the proof.

$$D_{q,t} \left( D_{q,y'_q} f \right) = 0.$$

**Remark 1**

Using Theorem 3.1, for  $n \rightarrow 1$ , the q-Euler Lagrange Equation (3) is equivalent to:

$$\sum_{k=0}^{\infty} q^k \left\{ \frac{f(q^k t, y(q^k t), y'_q(q^k t)) - f(q^k t, y(q^{k+1} t), y'_q(q^k t))}{y(q^k t)} \right\} = \frac{l(t)}{t}.$$

**Examples**

Recall that, the classical standard example, for  $f$  given by:

$$f(t, y, y') = \sqrt{1 + y'^2}$$

We get:

$$y = At + B$$

That is, the function must have constant first derivative and thus its graph is a straight line.

Now, we will study the q-analogue of this standard example.

**Theorem 4.1**

Let  $f$  given by:

$$f(t, y_q, y'_q) = \sqrt{1 + y_q'^2}$$

Satisfying the q-Euler Lagrange equation:

$$D_{q,y} f - D_{q,t} \left( D_{q,y'_q} f \right) = 0. \tag{7}$$

Then, we get  $y_q = At + C$ , for constants A and C.

**Proof**

Using equation (5), we obtain:

$$D_{q,y'_q} f = \frac{\sqrt{1 + y_q'^2} - \sqrt{1 + q y_q'^2}}{y_q' (1 - q)}.$$

Since, we have:

$$D_{q,y} f = 0$$

Then, by Equation (7), we get:

Therefore, we obtain:

$$D_{q,y} f = c,$$

Then, we get:

$$\frac{\sqrt{1 - y_q'^2} - \sqrt{1 + q y_q'^2}}{y_q' (1 - q)} = c.$$

This gives:

$$\begin{aligned} \sqrt{1 + y_q'^2} - \sqrt{1 + q y_q'^2} &= c y_q' (1 - q) \\ \sqrt{1 + q y_q'^2} - \sqrt{1 + q^2 y_q'^2} &= c q y_q' (1 - q) \\ &\vdots \\ \sqrt{1 + q^{n-1} y_q'^2} - \sqrt{1 + q^n y_q'^2} &= c q^{n-1} y_q' (1 - q). \end{aligned}$$

From which we get:

$$\sqrt{1 + y_q'^2} - \sqrt{1 + q^n y_q'^2} = c q^{n-1} y_q' (1 - q) + \dots + c y_q' (1 - q).$$

Then, we deduce:

$$\begin{aligned} \sqrt{1 + y_q'^2} - \sqrt{1 + q^n y_q'^2} &= c(1 - q) y_q' \sum_{k=0}^{n-1} q^k \\ &= c(1 - q) y_q' \left( \frac{1 - q^n}{1 - q} \right) \\ &= c [n]_q (1 - q) y_q'. \end{aligned}$$

As  $n \rightarrow \infty$ , we obtain:

$$\sqrt{1 + y_q'^2} - 1 = c y_q'$$

Which implies:

$$\sqrt{1 + y_q'^2} = 1 + c y_q'.$$

This gives:

$$1 + y_q'^2 = 1 + c^2 y_q'^2 + 2c y_q'.$$

Finally, we get:

$$y_q'^2 (1 - c^2) = 2c y_q'$$

Which gives:

$$y'_q = 0 \text{ or } y'_q = \frac{2c}{1-c^2}.$$

Let  $\frac{2c}{1-c^2} = A$ . Then  $y_q = Ax + C$ , this completes the proof.

Using Euler-Lagrange Equation (2) to find the extremals for the following function:

$$x \left( \frac{\partial}{\partial x} y(x) \right) + \left( \frac{\partial}{\partial x} y(x) \right)^2.$$

We obtain the following solution:

$$y(x) = \frac{-1}{4}x^2 + c_1x + c_2 \tag{8}$$

As q-deformation of this example, we get the following theorem.

**Theorem 4.2**

Let  $f$  given by:

$$f(t, y_q, y'_q) = ty'_q + y_q'^2$$

Satisfying the q-Euler Lagrange equation:

$$D_{q,y}f - D_{q,t}(D_{q,y'_q}f) = 0. \tag{9}$$

Then, we get:

$$y(t) = -\frac{1}{[2]_q^2}t^2 + c_1t + c_2.$$

**Proof**

Using Equation 5, we obtain:

$$\begin{aligned} D_{q,y'_q}f(t, y_q, y'_q) &= \frac{f(t, y_q, y'_q) - f(t, y_q, qy'_q)}{y'_q(1-q)} \\ &= \frac{(ty'_q + y_q'^2) - (tqy'_q + q^2y_q'^2)}{y'_q(1-q)} \\ &= \left( \frac{1}{1-q} \right) (t + y'_q - qt - q^2y'_q) \\ &= t + y'_q \left( \frac{1-q^2}{1-q} \right) \\ &= t + (1+q)y'_q = t + [2]_q y'_q. \end{aligned}$$

Since, we have:

$$D_{q,y}f(t, y_q, y'_q) = 0$$

Then:

$$D_{q,y} \left( D_{q,y'_q} f(t, y, y'_q) \right) = 0$$

Therefore, we get:

$$D_{q,y'_q} f(t, y, y'_q) = c$$

This gives:

$$t + [2]_q y'_q = c$$

Which is equivalent to:

$$t + [2]_q \frac{y(t) - y(qt)}{t(1-q)} = c.$$

This implies that:

$$[2]_q \left( \frac{y(t) - y(qt)}{t(1-q)} \right) = c - t.$$

Then, we get:

$$\begin{aligned} y(t) - y(qt) &= \frac{(c-t)t(1-q)}{[2]_q} = (-t^2 + ct) \left( \frac{1-q}{[2]_q} \right) \\ y(qt) - y(q^2t) &= (-q^2t^2 - cqt) \left( \frac{1-q}{[2]_q} \right) \\ &\vdots \\ y(q^{k-1}t) - y(q^kt) &= (-q^{2(k-1)}t^2 - cq^{(k+1)}t) \left( \frac{1-q}{[2]_q} \right), \end{aligned}$$

We deduce:

$$y(t) - y(q^kt) = \sum_{i=0}^{k-1} (-q^{2i}t^2 + cq^i t) \left( \frac{1-q}{[2]_q} \right).$$

But:

$$\begin{aligned} \sum_{i=0}^{k-1} -q^{2i}t^2 &= -t^2 \sum_{i=0}^{k-1} (-q^i)^i \\ &= -t^2 \left( \frac{1-q^{2k}}{1-q^2} \right) \end{aligned}$$

In addition:

$$\begin{aligned} \sum_{i=0}^{k-1} cq^i t &= ct \sum_{i=0}^{k-1} q^i \\ &= ct \left( \frac{1-q^k}{1-q} \right). \end{aligned}$$

Then, we get:

$$y(t) = y(q^k t) - \frac{(1-q)}{[2]_q} \left( \frac{1-q^{2k}}{1-q^2} \right) t^2 + c \frac{(1-q)}{[2]_q} \left( \frac{1-q^k}{1-q} \right) t$$

$$= y(q^k t) - \frac{(1-q^{2k})}{[2]_q^2} t^2 + c \frac{(1-q^k)}{[2]_q} t.$$

As  $k \rightarrow \infty$ , we obtain:

$$y(t) = y(0) - \frac{1}{[2]_q^2} t^2 + \frac{c}{[2]_q} t$$

$$= -\frac{1}{[2]_q^2} t^2 + c_1 t + c_2$$

where,  $c_1 = \frac{c}{[2]_q}$  and  $c_2 = y(0)$ .

### Remark 2

As  $q \rightarrow \infty$ , Equation (10) becomes:

$$y(t) = \frac{-1}{2^2} t^2 + c_1 t + c_2$$

$$= -\frac{1}{4} t^2 + c_1 t + c_2$$

Which gives the classical case studied see Equation (8).

### Author's Contributions

All authors equally contributed in this work.

### Ethics

This article is original and contains unpublished material. The corresponding author confirms that all of the other authors have read and approved the manuscript and no ethical issues involved.

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